

KO'P QATLAMLI STRUKTURALARDA TOK TASHUVCHILARNING KINETIKASINING NAZARIY TADQIQI. 2-QISM

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Annotatsiya: *Shredinger tenglamasi asosida markazida potensial to'siq mavjud bo'lgan besh qavatli yarimo'tkazgich nanostruktura uchun nazariy hisoblashlar ishlab chiqildi. Qo'zg'alishlar nazariyasiga ko'ra asosiy va qo'zg'atilgan holat energiyalarning o'rta potensial o'ra balandligi hamda uchinchi va ikkinchi qatlamlar o'lchamlarining nisbatiga bog'lanishilari hisoblangan va ular elektronlar energiyasining qatlamlar o'lchamlariga holda siljishiga oloib kelishi va bu siljishning musbat yoki manfiy bo'lishi qaralayotgan energetiya sohalariga, qatlamlarning zonaviy parametrlariga bog'liqligi ko'rsatib o'tilgan.*

Kalit so'zlar: *potensial o'ra, potensial to'siq, elektron, energetik spektr, qo'zg'alishlar nazariyasi, effektiv massa, qatlamli struktura.*

ТЕОРЕТИЧЕСКОЕ ИССЛЕДОВАНИЕ КИНЕТИКИ НОСИТЕЛЕЙ ТОКА В МНОГОСЛОЙНЫХ СТРУКТУРАХ. ЧАСТЬ 2

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Аннотация: На основе уравнения Шрёдингера разработана теоретическая модель пятислойной полупроводниковой наноструктуры с потенциальной ямой в центральной области.

Согласно теории возмущений рассчитаны энергии основного и возбужденного состояний в зависимости от высоты потенциального барьера, находящегося в центре структуры и отношением размеров третьего и второго слоев и показано, что они приводят к сдвигу энергии электронов, зависящему от размеров слоев. Причем положительное или отрицательное значение этого сдвига зависит от зонных параметров, а также от размеров слоев и зонных параметров слоев.

Ключевые слова: потенциальная яма, потенциальный барьер, электроны, энергетический спектр, теория возмущений, эффективная масса, слоистая структура.

THEORETICAL STUDY OF THE KINETICS OF CURRENT CARRIERS IN MULTILAYER STRUCTURES. PART 1

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Abstract: *A theoretical model of a five-layer semiconductor nanostructure with a potential well in the central region has been developed using the Schrödinger equation.*

Using perturbation theory, the energies of the ground and excited states have been calculated as functions of the height of the potential barrier located in the center of the structure and the ratio of the sizes of the third and second layers. It has been shown that these values lead to a shift in the electron energy that depends on the layer sizes. Furthermore, the positive or negative value of this shift depends on the band parameters, as well as the layer sizes and band parameters of the layers.

Keywords: *potential well, potential barrier, electrons, energy spectrum, perturbation theory, effective mass, layered structure.*

Kirish

Qator adabiyotlarda ko‘p qatlamli geteroo‘tikli strukturalarda elektronli holatlar chekli energiyaviy chuqurlikli to‘g‘ri burchakli ko‘rinishdagi potensial o‘radagi tok tashuvchilarning harakati va tunellanishi masalasi kabi oddiy yaqinlashishda hal qilingan [1-11].

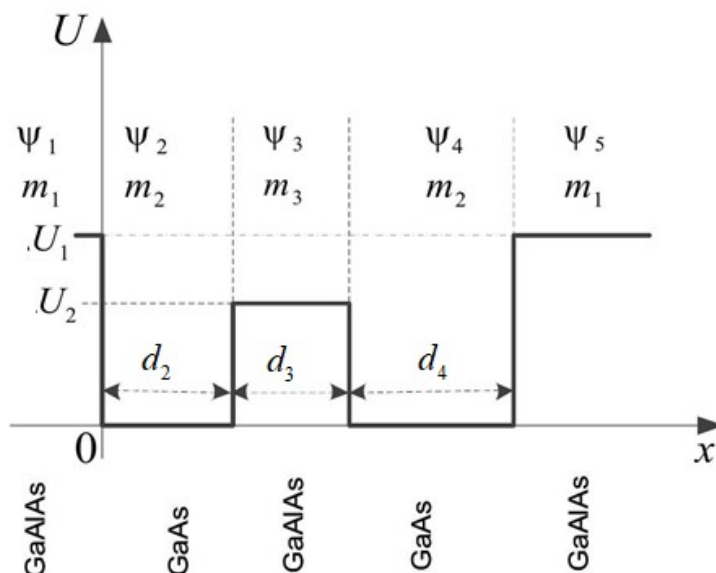
Adabiyotlar tahlili

Nanoelektronikaning rivojlanishi hisobiga bunday masalalarga qayta murojaat qilinib, murakkab ko‘rinishdagi qator potentsiallar uchun bunday masalalar yechilgan [1-6]. Ammo ixtiyoriy ko‘rinishdagi potentsialdagi mikrozarraning harakatiga tegishli kvant mexanikaviy masalalar kvaziklassik yaqinlashishda yechishni talab etadi. Shu sababli quyida nanostrukturalar, jumladan geteroo‘tikli tizimlar, o‘ta panjaralardagi

tok tashuvchilarning energetik spektrini hisoblashda qo'llaniladigan kvaziklassik hisoblash usuliga to'xtalib o'tamiz.

Tadqiqot metodologiyasi

Ayni paytda rezonans tunnelli diod (RTS) nanoelektronikaning istiqbolli qurilmalarining asosini tashkil etadi. Xususan, RTSlarining yangi avlodlarini, masalan, elektron va optik cheklanishi ajratilgan geterolazerlarni yaratish uchun [7, 8] markazida qo'shimcha to'g'ri burchakli potensial o'raga ega bo'lgan ko'p qatlamli strukturalardan foydalaniladi (1-rasm). Energetik chuqurlik tufayli potensial o'raning keng asosiy qismida kvant sathlarining joylashuvini boshqarish uchun yangi imkoniyatlar yuzaga keladi.



1-rasm. Rezonansli bog'langan kvant potensial o'ralar sistemasi.

Ushbu rasmda m_i - zaryad tashuvchilarning effektiv massalari; U_1 va U_2 - potensial o'ralarning balandliklari; d_2, d_3, d_4 - struktura ichki qatlamlar qalinliklari. Shunday qilib, 1-rasmda qayd qilingan potensial maydonda elektronlarning hatti harakati bu potensialni o'tkazuvchanlik zonasi ichiga olgan Shredenger tenglamasini yechish bilan hal etiladi. Har bir qatlam uchun Shredinger tenglamasining ($E < U_1, U_2$ da) yechimi quyidagi ko'rinishda ifodalanadi:

$$\psi_1 = A_1 \cdot \exp(k_1 x), \psi_2 = A_2 \cdot \cos(k_2 x) + B_2 \cdot \sin(k_2 x), \psi_3 = A_3 \cdot \operatorname{ch}(k_3 (x - d_3)) + B_3 \cdot \operatorname{sh}(k_3 (x - d_3)),$$

$$\psi_4 = A_4 \cdot \cos(k_2 (x - d_2 - d_3)) + B_4 \cdot \sin(k_2 (x - d_1 + d_2)), \psi_5 = B_5 \cdot \exp(-k_1 (x - d_1 - d_2 - d_3))$$

bu yerda $k_1 = \sqrt{2m_1 \hbar^{-2} |U_1 - E|}, k_2 = \sqrt{2m_2 \hbar^{-2} E}, k_3 = \sqrt{2m_3 \hbar^{-2} |U_2 - E|}$.

Bastard sharti [1] e'tiborga olinb, A_i, B_i noma'lum noma'lumlar oldidagi

koeffitsiyentlardan tuzilgan determinant nolga teng bo'lish shartidan holatlar energiyasini aniqlash uchun quyidagi ifodani olamiz:

$$\begin{aligned} & \left[-\tilde{k}_2 (\tilde{k}_1 - \tilde{k}_3) \cos(k_2 d_2) + (\tilde{k}_1 \tilde{k}_3 + \tilde{k}_2^2) \sin(k_2 d_2) \right] \left[\tilde{k}_4 \cos(k_4 d_4) + \tilde{k}_5 \sin(k_4 d_4) \right] \times \\ & \times \left[\tilde{k}_4 \operatorname{sh}(k_3 d_3) - \tilde{k}_3 \operatorname{ch}(k_3 d_3) \right] = 0. \end{aligned} \quad (1)$$

Xususan $U_1 \geq E > U_2$ shartda:

$$\begin{aligned} & \left[\tilde{k}_2^2 (\tilde{k}_3^2 - \tilde{k}_1^2) + (\tilde{k}_1^2 \tilde{k}_3^2 - \tilde{k}_2^4) \operatorname{tg}(k_2 d_4) \right] \operatorname{tg}(k_3 d_3) + 2\tilde{k}_1 \tilde{k}_2 \tilde{k}_3 \cdot \left[\operatorname{tg}(k_2 d_2) \operatorname{tg}(k_2 d_4) - 1 \right] + \\ & + \tilde{k}_2 \left[\tilde{k}_3 (\tilde{k}_2^2 - \tilde{k}_1^2) + \tilde{k}_1 (\tilde{k}_3^2 + \tilde{k}_2^2) \operatorname{tg}(k_3 d_3) \right] \left[\operatorname{tg}(k_2 d_2) + \operatorname{tg}(k_2 d_4) \right] = 0, \end{aligned} \quad (2)$$

bu yerda $\tilde{k}_{n'n} = \tilde{k}_{n'}/\tilde{k}_n = m_n k_{n'}/k_n m_{n'}$. Potensial to'siqi balandligi $U_1 (U_1 \gg U_2)$ shartni

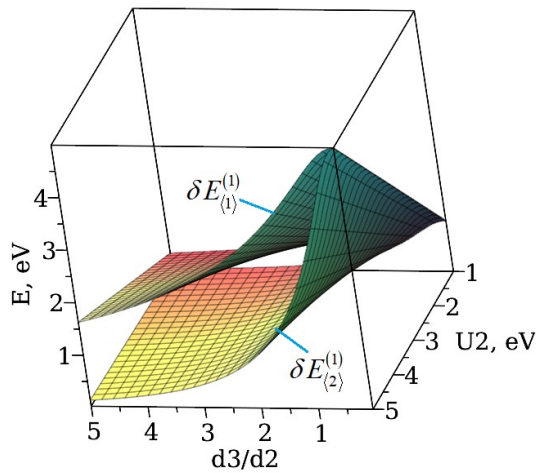
qanoatlantirganda $E < U_2$ va $d_2 = d_4$ shartda: $\left[\tilde{k}_2^2 + \tilde{k}_3^2 \operatorname{tg}(k_2 d_2) \operatorname{tg}(k_2 d_4) \right] \operatorname{th}(k_3 d_3) +$
 $+ \tilde{k}_2 \tilde{k}_3 \left[\operatorname{tg}(k_2 d_2) + \operatorname{tg}(k_2 d_4) \right] = 0$; $E > U_2$ shartda esa: $\left[\tilde{k}_2^2 - \tilde{k}_3^2 \operatorname{tg}(k_2 d_2) \operatorname{tg}(k_2 d_4) \right] \operatorname{tg}(k_3 d_3) +$
 $+ \tilde{k}_2 \tilde{k}_3 \left[\operatorname{tg}(k_2 d_2) + \operatorname{tg}(k_2 d_4) \right] = 0$.

Natijalar va ularning tahlili

G'alayonlar nazariyasiga ko'ra asosiy $(n=1)$ va qo'zg'algan $(n=2)$ holat energiyalari cheksiz chuqur potensial o'ra yaqinlashishia qo'zg'alishlar nazariyasiga ko'ra:

$$\delta E_n^{(1)} = \left[\frac{\pi n^2 \hbar^2}{2(d_2 + d_3)^2} \left(\frac{1}{m_3} - \frac{1}{m_2} \right) - U_2 \right] \cdot \left[\frac{d_2}{d_2 + d_3} - \frac{(-1)^n}{n\pi} \sin \left(\frac{2\pi n d_2}{d_2 + d_3} \right) \right]. \quad (3)$$

(3) ifodaga asoslanib qo'zg'alishlar nazariyasiga ko'ra hisoblangan $\delta E_1^{(1)}$ asosiy va $\delta E_2^{(1)}$ qo'zg'atilgan holat energiyalarning o'rta potensial o'ra balandligi (U_2) hamda uchinchi va ikkinchi qatlamlar o'lchamlarining nisbati (d_3/d_2) ga bog'lanishlaridan ko'rinyaptiki, ular elektronlar energiyasining qatlamlar o'lchamlariga holda siljishiga oloib keladi va bu siljishning musbat yoki manfiy bo'lishi qaralayotgan energetik sohalariga, qatlamlarning zonaviy parametrlariga bog'liqdir.



2-rasm. Asosiy $\delta E_1^{(1)}$ va qo'zg'algan $\delta E_2^{(1)}$ holat energiyalarning o'rta potensial o'ra balandligi hamda uchinchi va ikkinchi qatlamlar o'lchamlarining nisbatiga bog'lanishi.

Sonli hisoblashlarda kattaliklarning quyidagi qiymatlari tanlangan [7, 8]:

$$d_2 = d_3 = d_4 = 20 \text{ \AA}, \quad U_1 = U_2 = 1,32 \text{ eV}, \\ m_1 = m_3 = 0,123m_0, \quad m_2 = 0,046m_0.$$

Xulosa

Shunday qilib, keltirilgan ifodalar va hisoblangan grafik bog'liqliklar shunday xulosa qilish imkonini beradiki, murakkab shaklli kvant potensial o'ralardagi energetik sathlarning joylashuvini sifat jihatdan tavsiflash uchun cheksiz chuqur o'ra modelidan foydalanish faqat nisbatan katta

chuqurlikka ega o'ralar uchungina mumkin. Murakkab shaklli kvant potensial o'ralarda energetik spektrni aniqlovchi ifodalar zamonaviy yuqori tezkor elektronika qurilmalarida qo'llaniladigan tunnellash–rezonansli tuzilmalarni loyihalash va modellashtirishda foydali bo'lishi mumkin.

Foydalanilgan adabiyotlar

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